Lecture 8
Using Entropy
Entropy – How to get that?

\[ dU = dq + dw \quad \text{First law} \]
\[ = dq_{rev} + dw_{rev} \]
\[ dw_{rev} = -PdV \]
\[ TdS = dq_{rev} \]
\[ dU = TdS - PdV \]

For an ideal gas,
\[ dU = nC_v \, dT \]
\[ P = nRT/V \]
\[ nC_v \, dT = TdS - (nRT/V) \, dV \]
\[ dS = nC_v \, dT/T + nR \, dV/V \]
If $C_v$ is constant for the temperature range,

$$\int_1^2 dS = n \ C_v \int_1^2 dT/T + nR \int_1^2 dV/V$$

$$\Delta S = n \ C_v \ln \frac{T_2}{T_1} + nR \ln \frac{V_2}{V_1}$$

In terms of $P$ and $T$,

$$V_2 = \frac{RT_2}{P_2}; \ V_1 = \frac{RT_1}{P_1}$$

$$\frac{V_2}{V_1} = \frac{RT_2}{P_2} \frac{P_1}{RT_1} = \frac{T_2 P_1}{T_1 P_2}$$

$$\Delta S = n \ C_v \ln \frac{T_2}{T_1} + nR \ln \frac{T_2}{T_1} - nR \ln \frac{P_2}{P_1}$$

Since $V_2/V_1 = T_2 P_1 / P_2 T_1$

$$= n \ (C_v + R) \ln \frac{T_2}{T_1} - nR \ln \frac{P_2}{P_1}$$

$$= n \ C_p \ln \frac{T_2}{T_1} - nR \ln \frac{P_2}{P_1}$$

Entropy has units cal/deg which is referred as the entropy unit.
For an isothermal process,
\[ \Delta S_T = nR \ln \frac{V_2}{V_1} = -nR \ln \frac{P_2}{P_1} \]

For isochoric process,
\[ \Delta S_V = nC_v \ln \frac{T_2}{T_1} \]

For isobaric process
\[ \Delta S_p = nC_p \ln \frac{T_2}{T_1} \]
A relation can be obtained in terms of \( P \) and \( V \) also.

\[
\Delta S = nC_v \ln \frac{T_2}{T_1} + nR \ln \frac{V_2}{V_1}
\]

\[
\frac{T_2}{T_1} = \frac{P_2 V_2}{P_1 V_1}
\]

\[
= nC_v \ln \frac{P_2}{P_1} + nC_v \ln \frac{V_2}{V_1} + nR \ln \frac{V_2}{V_1}
\]

\[
= nC_v \ln \frac{P_2}{P_1} + nC_p \ln \frac{V_2}{V_1}
\]
Entropy change for a phase change

$$TdS = dU + PdV$$

At constant pressure,

$$TdS = d(U + PV) = dH$$

$$dS = dH/T$$

$$\Delta S = \Delta H/T$$

Depending on the process, $$\Delta H = L_v, L_f, \text{ etc.}$$
Variation of enthalpy with temperature

\[ dS = \frac{dq_{rev}}{T} = \frac{dH}{T} = C_p \frac{dT}{T} \]
constant pressure

\[ dS = \frac{dq_{rev}}{T} = \frac{dU}{T} = C_v \frac{dT}{T} \]
constant volume

These equations can be integrated to get variation in entropy. If heat capacities are expressed as a function of temperature, \( \Delta S \) can be evaluated accurately.
Variation with temperature

Example

Let us say a solid is heated to melt and the liquid is made to boil.

\[ S(T) = S(0) + \int_{T_0}^{T} (C_p(s) dT/T) + \left( \Delta_{\text{fus}} H/T_f \right) + \int_{T_b}^{T_f} (C_p(l) dT/T) + \left( \Delta_{\text{eva}} H/T_b \right) + \int_{T_b}^{T_f} C_p(g) dT/T \]
$$\Delta S = \int_{T_1}^{T_2} \frac{C}{T} \, dT$$
Entropy changes with specific processes

\[ \Delta S = \Delta H / T. \] When \( \Delta H \) is positive such as melting, \( \Delta S \) is positive.

Standard molar entropy of vaporization of liquids is a constant of value, 85 J/K/mol. This is called the Trouton’s rule.

With this rule, enthalpy of vaporization of liquids at their normal boiling points can be calculated. When liquids are associated, the entropies are generally larger.
## Entropies of phase transitions (Table)

Entropies of phase transitions at 1 atm ($\Delta S_m$/JK$^{-1}$mol$^{-1}$)

<table>
<thead>
<tr>
<th>System</th>
<th>Fusion</th>
<th>Evaporation</th>
</tr>
</thead>
<tbody>
<tr>
<td>He</td>
<td>6.3</td>
<td>19.7</td>
</tr>
<tr>
<td>Ar</td>
<td>14.2</td>
<td>71.7</td>
</tr>
<tr>
<td>H$_2$</td>
<td>8.4</td>
<td>44.4</td>
</tr>
<tr>
<td>N$_2$</td>
<td>11.42</td>
<td>73.16</td>
</tr>
<tr>
<td>H$_2$O</td>
<td>22.0</td>
<td>109.0</td>
</tr>
<tr>
<td>NH$_3$</td>
<td>28.9</td>
<td>97.21</td>
</tr>
<tr>
<td>CH$_3$OH</td>
<td>18.1</td>
<td>104</td>
</tr>
<tr>
<td>CH$_3$COOH</td>
<td>40.4</td>
<td>61.9</td>
</tr>
</tbody>
</table>
Entropy temperature diagram
Heat absorbed \( q_2 = T_2 (S_2 - S_1) \)

Heat evolved during isothermal compression
\[-q_1 = T_1 (S_1 - S_2)\]
\[= -T_1 (S_2 - S_1)\]

This value of energy is the energy unavailable for work.

Heat absorbed \( = \) work done
\[= q_2 - q_1\]
\[= (T_2 - T_1) (S_2 - S_1)\]

This product is the area \( \text{ABCD} \)
The heat energy transferred cannot be completely used for work. **Energy unavailable for work is proportional to \( \Delta S \).**

**Maximum work**

\[
W_{\text{max}} = \text{Efficiency} \times \text{heat absorbed} \\
= \left( \frac{(T_2 - T_1)}{T_2} \right) \times q_2 \\
= q_2 \left( 1 - \frac{T_1}{T_2} \right) = q_2 - T_1 \frac{q_2}{T_2} \\
= q_2 - T_1 \Delta S
\]
Using Thermodynamic Equations
\[ \text{d}U = T \text{d}S - P \text{d}V \quad (1) \quad \textit{Fundamental equation} \]

Applicable to both reversible and irreversible.

\[ U = f(S,V) \]

\[ \text{d}U = \left( \frac{\partial U}{\partial S} \right)_V \text{d}S + \left( \frac{\partial U}{\partial V} \right)_S \text{d}V \quad (2) \]

Compare (1) and (2)

\[ \left( \frac{\partial U}{\partial S} \right)_V = T \quad (3) \]

\[ \left( \frac{\partial U}{\partial V} \right)_S = -P \quad (4) \text{ Thermodynamics give unusual relations!} \]

From (3), \( \frac{\partial U}{\partial V} \frac{\partial S}{\partial V} = \left( \frac{\partial T}{\partial V} \right)_S \)

From (4), \( \frac{\partial U}{\partial S} \frac{\partial V}{\partial V} = -\left( \frac{\partial P}{\partial S} \right)_V \)

As \( \text{d}U \) is an exact differential,

\[ \left( \frac{\partial T}{\partial V} \right)_S = -\left( \frac{\partial P}{\partial S} \right)_V \quad \text{Maxwell’s relation (1)} \]
These are four fundamental relations of thermodynamics. These are actually several ways of looking at the same relation, $dU = TdS - PdV$.

Each relation relates one thermodynamic function to two variables. S and V are said to be natural variables of U; S and P for H; T and V for A as well as T and P for G.
Maxwell’s relations

\[ \frac{\partial T}{\partial V} \bigg|_S = -\frac{\partial P}{\partial S} \bigg|_V \]
\[ \frac{\partial T}{\partial P} \bigg|_S = \frac{\partial V}{\partial S} \bigg|_P \]
\[ \frac{\partial S}{\partial V} \bigg|_T = -\frac{\partial P}{\partial T} \bigg|_V \]
\[ \frac{\partial S}{\partial P} \bigg|_T = \frac{\partial V}{\partial T} \bigg|_P \]

The first two relations are for constant entropy, i.e., adiabatic processes. Rate of change of temperature with volume and pressure are the quantities discussed. The later two relations are for constant temperature, i.e., isothermal processes. Entropy change as a function of volume and pressure can be related to changes in pressure and volume.

How can such relations be tested?
Thermodynamic equation of state

The ideal gas law and van der Waals equation are relations between $P$, $V$ and $T$. These are based on some data and extrapolations and speculations on the molecular dimensions and interactions between molecules.

Is there a more generalised law describing equilibrium?

Thermodynamics gives one.
dU = TdS – PdV is the condition of equilibrium.

Let us write,

\[(\partial U)_T = T(\partial S)_T - P(\partial V)_T\]

Divide by \((\partial V)_T\)

\[\frac{(\partial U)}{(\partial V)}_T = T\frac{(\partial S)}{(\partial V)}_T - P\]

U and S are functions of T and V.
Thus, the above relation relates P with functions of T and V. Therefore, this is an equation of state.
Using value of \((\partial S/\partial V)_T\) from Maxwell’s relation \((\partial S/\partial V)_T = (\partial P/\partial T)_V\),

\[ P = T(\partial P/\partial T)_V - (\partial U/\partial V)_T \]

Taking,

\[ dH = TdS + VdP \]

\[ (\partial H/\partial P)_T = T(\partial S/\partial P)_T + V \]

Using Maxwell relation, \(-(\partial S/\partial P)_T = (\partial V/\partial T)_P\)

\[ V = T(\partial V/\partial T)_P + (\partial H/\partial P)_T \]

This equation expresses volume as a function of temperature and pressure.

These relations are applicable to any substance.